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**Reliable Solution of a Unilateral
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Technical report No. 860

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Reliable Solution of a Unilateral Frictionless Contact Problem in Quasi-Coupled Thermo-Elasticity with Uncertain Input Data

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Abstract:

A unilateral contact problem without friction in quasi-coupled thermo-elasticity and with uncertain input data is analysed. The worst scenario method is used to find the most "dangerous" admissible input data.

Keywords:

Worst scenario method, contact problem, uncertain data, thermo-elasticity

1 Introduction

In this contribution we deal with contact problems without friction (see [4], [5], [6]) in quasi-coupled thermo-elasticity considering uncertain input data representing extension of problems solved in [5] and [6]. By uncertain input data we mean physical coefficients, right-hand sides, etc., which cannot be determined uniquely but only in some intervals determined by the measurements. The reliable solution is defined as the worst among a set of possible solutions, and the degree of badness is measured by a criterion-functional (see [1]). The main aim of our contribution will be to find maximal values of this functional. We prove the solvability of the corresponding maximization (worst scenario) problems.

2 Formulation of the Problem

Let us assume a union Ω of bounded domains Ω^ι , $\iota = 1, \dots, s$, with Lipschitz boundaries $\partial\Omega^\iota$, occupied by elastic bodies such that $\Omega = \bigcup_{\iota=1}^s \Omega^\iota \subset R^2$. Let the boundary $\partial\Omega = \bigcup_{\iota=1}^s \partial\Omega^\iota$ consist of three disjoint parts Γ_τ , Γ_u and Γ_c , such that $\partial\Omega = \bar{\Gamma}_\tau \cup \bar{\Gamma}_u \cup \bar{\Gamma}_c$, $\Gamma_c = \bigcup_{k,l} \Gamma^{kl}$, $\Gamma^{kl} = \partial\Omega^k \cap \partial\Omega^l$, $1 \leq k, l \leq s$, for $k \neq l$, and $\bar{\Gamma}_\tau, \bar{\Gamma}_u, \bar{\Gamma}_c$ denotes the closures in $\partial\Omega$.

Let the heat sources W^ι , the prescribed temperature T_1 , the body forces \mathbf{F} , the surface forces \mathbf{P} , displacements \mathbf{u}_0 , elastic coefficients c_{ijkl} , coefficients of thermal expansion β_{ij} and the reference temperature T_0 be given. Throughout the paper we use the summation convention, i.e. a repeated index implies summation from 1 to 2. Furthermore, $\mathbf{n}^k = (n_i^k)$, $i = 1, 2, 1 \leq k \leq s$, denotes the unit normal with respect to $\partial\Omega^k$, $\mathbf{n}^k = -\mathbf{n}^l$ on Γ^{kl} . Assume that κ^ι and C^ι are positive definite symmetric matrix functions,

$$\begin{aligned} 0 < \kappa_0^\iota \leq \kappa_{ij}^\iota \zeta_i \zeta_j |\zeta|^{-2} \leq \kappa_1^\iota < +\infty \quad \text{for a.a. } \mathbf{x} \in \Omega^\iota, \zeta \in R^2, \\ 0 < c_0^\iota \leq c_{ijkl}^\iota \xi_{ij} \xi_{kl} |\xi|^{-2} \leq c_1^\iota < +\infty \quad \text{for a.a. } \mathbf{x} \in \Omega^\iota, \xi \in R^4, \xi_{ij} = \xi_{ij}, \end{aligned}$$

where $\kappa_0^\iota, \kappa_1^\iota, c_0^\iota, c_1^\iota$ are constants independent of $\mathbf{x} \in \Omega^\iota$. Let $\kappa_{ij}^\iota \in L^\infty(\Omega^\iota)$, $W^\iota \in L^2(\Omega^\iota)$, $T_1^\iota \in H^1(\Omega^\iota)$, $T_1^k = T_1^l$ on $\bigcup_{k,l} \Gamma^{kl}$, $c_{ijkl}^\iota \in L^\infty(\Omega^\iota)$, $F_i^\iota \in L^2(\Omega^\iota)$, $P_i \in L^2(\Gamma_\tau)$, $\beta_{ij}^\iota \in L^\infty(\Omega^\iota)$, $\mathbf{u}_0 \in [H^1(\Omega^\iota)]^2$.

We will deal with the following problem:

Problem (P): Find a pair of functions (T, \mathbf{u}) satisfying

$$\frac{\partial}{\partial x_i} \left(\kappa_{ij}^\iota \frac{\partial T^\iota}{\partial x_j} \right) + W^\iota = 0, \quad \frac{\partial}{\partial x_j} \tau_{ij}(\mathbf{u}^\iota, T^\iota) + F_i^\iota = 0 \quad \text{in } \Omega^\iota, \quad 1 \leq \iota \leq s, \quad i = 1, 2 \quad (2.1)$$

$$\tau_{ij}(\mathbf{u}^\iota, T^\iota) = c_{ijkl}^\iota e_{kl}(\mathbf{u}^\iota) - \beta_{ij}^\iota (T^\iota - T_0^\iota) \quad \text{in } \Omega^\iota, \quad 1 \leq \iota \leq s, \quad i = 1, 2 \quad (2.2)$$

$$\kappa_{ij}^\iota \frac{\partial T}{\partial x_j} n_i = 0, \quad \mathbf{u} = \mathbf{u}_0 \quad \text{on } \Gamma_u, \quad (2.3)$$

$$T = T_1, \quad \tau_{ij}(\mathbf{u}, T) n_j = P_i \quad \text{on } \Gamma_\tau, \quad (2.4)$$

$$T^k = T^l, \quad \left(\kappa_{ij} \frac{\partial T}{\partial x_j} n_i \right)^k - \left(\kappa_{ij} \frac{\partial T}{\partial x_j} n_i \right)^l = 0 \quad \text{on } \bigcup_{k,l} \Gamma^{kl}, \quad 1 \leq k, l \leq s, \quad (2.5)$$

$$u_n^k - u_n^l \leq 0, \quad \tau_n^k \leq 0, \quad (u_n^k - u_n^l) \tau_n^k = 0 \quad \text{on } \bigcup_{k,l} \Gamma^{kl}, \quad 1 \leq k, l \leq s, \quad (2.6)$$

$$\tau_t^k = -\tau_t^l = 0 \quad \text{on } \bigcup_{k,l} \Gamma^{kl}, \quad 1 \leq k, l \leq s, \quad (2.7)$$

where $e_{ij}(\mathbf{u}) = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$, $u_n^k = u_i^k n_i^k$, $u_n^l = u_i^l n_i^l = -u_i^k n_i^k$ (no sum over k or l), $u_t^k = (u_{ti}^k)$,

$u_{ti}^k = u_i^k - u_n^k n_i^k$, $u_t^l = (u_{ti}^l)$, $u_{ti}^l = u_i^l - u_n^l n_i^l$, $i = 1, 2$, $\tau_n^k = \tau_{ij}^k n_i^k n_j^k$, $\tau_t^k = (\tau_{ti}^k)$, $\tau_{ti}^k = \tau_{ij}^k n_j^k - \tau_n^k n_i^k$, $\tau_n^l = \tau_{ij}^l n_i^l n_j^l$, $\tau_t^l = (\tau_{ti}^l)$, $\tau_{ti}^l = \tau_{ij}^l n_j^l - \tau_n^l n_i^l$.

Since the stress and strain tensors and coefficient of thermal expansion are symmetric then the entries of any symmetric 3×3 matrices $\{\tau_{ij}\}$ can be rewritten in the vector notation $\{\tau_j\}$, $j = 1, 2, 3$ and similarly the symmetric matrices $\{e_{ij}\}$, $\{\beta_{ij}\}$ by vectors $\{e_j\}$, $\{\beta_j\}$. Then (2.2) can be rewritten as

$$\tau_i(\mathbf{u}^t, T^t) = \sum_{j=1}^3 A_{ij}^t e_j(\mathbf{u}^t) - \beta_i^t (T^t - T_0^t) \quad \text{in } \Omega^t, \quad 1 \leq t \leq s, \quad 1 \leq i, j \leq 3, \quad (2.8)$$

where A^t is a symmetric 3×3 matrix, $A_{ik}^t \in L^\infty(\Omega^t)$, $i = 1, \dots, s$. Since $\tau_{ij} e_{ij} = \sum_{i=1}^2 \tau_i e_i + 2\tau_3 e_3$, we can write

$$c_{ijkl}^t e_{ij} e_{kl} = \sum_{i,j=1}^3 B_{ij}^t e_i e_j,$$

where B^t is a symmetric 3×3 matrix such that $B_{ij}^t = A_{ij}^t$ for $i, j = 1, 2$, $B_{ij}^t = \frac{3}{2} A_{ij}^t$ for $i = 1, 2$, $j = 3$ and $B_{ij}^t = 2A_{ij}^t$ for $i, j = 3$.

In what follows, we denote

$$\begin{aligned} W_1 &= \prod_{i=1}^s H^1(\Omega^i), \quad \|w\|_{W_1} = \left(\sum_{i \leq s} \|w^i\|_{1, \Omega^i}^2 \right)^{\frac{1}{2}}, \\ W &= \prod_{i=1}^s [H^1(\Omega^i)]^2, \quad \|v\|_W = \left(\sum_{i \leq s} \sum_{i \leq 2} \|v_i^i\|_{1, \Omega^i}^2 \right)^{\frac{1}{2}}, \\ V_1 &= \left\{ z \mid z \in W_1, z = 0 \text{ on } \Gamma_\tau, z^k = z^l \text{ on } \bigcup_{k,l} \Gamma^{kl} \right\}, \\ V &= \{ \mathbf{v} \mid \mathbf{v} \in W, \mathbf{v} = 0 \text{ on } \Gamma_u \}, \quad K = \left\{ \mathbf{v} \mid \mathbf{v} \in V, v_n^k - v_n^l \leq 0 \text{ on } \bigcup_{k,l} \Gamma^{kl} \right\}. \end{aligned}$$

Definition 1. We say that the pair of functions T and \mathbf{u} is a weak solution of problem (P), if $T - T_1 \in V_1$,

$$b(T, z) = s(z) \quad \forall z \in V_1, \quad (2.9)$$

$$\mathbf{u} - \mathbf{u}_0 \in K,$$

$$a(\mathbf{u}, \mathbf{v} - \mathbf{u}) \geq S(\mathbf{v} - \mathbf{u}, T) \quad \forall \mathbf{v} \in \mathbf{u}_0 + K, \quad (2.10)$$

where

$$b(T, z) = \sum_{i=1}^s \int_{\Omega^i} \kappa_{ij}^i \frac{\partial T^i}{\partial x_i} \frac{\partial z^i}{\partial x_j} d\mathbf{x}, \quad s(z) = \sum_{i=1}^s \int_{\Omega^i} W^i z^i d\mathbf{x}, \quad (2.11)$$

$$a(\mathbf{u}, \mathbf{v}) = \sum_{i=1}^s \int_{\Omega^i} \sum_{i,j=1}^3 B_{ij}^i e_i(\mathbf{u}^i) e_j(\mathbf{v}^i) d\mathbf{x}, \quad (2.12)$$

$$S(\mathbf{v}, T) = \sum_{i=1}^s \int_{\Omega^i} F_i^i v_i^i d\mathbf{x} + \int_{\Gamma_\tau} P_i v_i ds - \sum_{i=1}^s \int_{\Omega^i} \beta_i^i (T^i - T_0^i) v_i^i d\mathbf{x}. \quad (2.13)$$

Remark 1. In $S(\mathbf{v}, T)$ we insert the weak solution T of (2.9). Moreover, we assume that \mathbf{u}_0 satisfies

$$u_{0n}^k - u_{0n}^l = 0 \quad \text{on} \quad \bigcup_{k,l} \Gamma^{kl}. \quad (2.14)$$

3 Worst Scenario Method for Uncertain Input Data

Let us assume that the input data

$$A = \{B^\iota, \kappa^\iota, F_i^\iota, W^\iota, \beta_i^\iota, P_i, u_{0i}, T_1, \iota = 1, \dots, s\}$$

are uncertain, and belong to some sets of admissible data, i.e.

$$\begin{aligned} A \in U_{ad} &\Leftrightarrow B^\iota \in U_{ad}^{B^\iota}, \kappa^\iota \in U_{ad}^{\kappa^\iota}, F_i^\iota \in U_{ad}^{F_i^\iota}, W^\iota \in U_{ad}^{W^\iota}, \beta_i^\iota \in U_{ad}^{\beta_i^\iota}, \\ P &\in U_{ad}^{P_i}, u_{0i} \in U_{ad}^{u_{0i}}, T_1 \in U_{ad}^{T_1}. \end{aligned}$$

We will assume that all the bodies Ω^ι are piecewise homogeneous, so that partitions of $\bar{\Omega}^\iota$ exist such that

$$\bar{\Omega}^\iota = \bigcup_{j=1}^{r^\iota} \bar{\Omega}_j^\iota, \quad \Omega_j^\iota \cap \Omega_k^\iota = \emptyset \quad \text{for } j \neq k, \quad 1 \leq \iota \leq s, \quad (3.1)$$

$$\Gamma^{kl} = \bigcup_{q=1}^{Q_{kl}} \bar{\Gamma}_q^{kl}, \quad \Gamma_q^{kl} \cap \Gamma_p^{kl} = \emptyset \quad \text{for } q \neq p, \quad \forall k, l. \quad (3.2)$$

Let the data $B^\iota, \kappa^\iota, F^\iota, W^\iota, \beta^\iota$ be piecewise constant with respect to the corresponding partitioning (3.1) and let us denote

$$\Gamma_u^\iota = \Gamma_u \cap \partial\Omega^\iota, \quad \iota = 1, \dots, s \quad \text{and} \quad \Gamma_\tau^\iota = \Gamma_\tau \cap \partial\Omega^\iota, \quad \iota \leq s. \quad (3.3)$$

Further, we define the sets of admissible matrices:

$$\begin{aligned} U_{ad}^{B^\iota} &= \{3 \times 3 \text{ symmetric matrices } B^\iota : \underline{B}_{ik}^\iota(j) \leq B_{ik}^\iota|_{\Omega_j^\iota} = \text{const.} \leq \bar{B}_{ik}^\iota(j), \\ &j \leq r^\iota, \quad i, k = 1, \dots, 3\} \end{aligned} \quad (3.4)$$

where $\underline{B}^\iota(j)$ and $\bar{B}^\iota(j)$ are given 3×3 symmetric matrices, $\iota = 1, \dots, s$, and let there exist positive constants $c_B^\iota(j)$ such that

$$\begin{aligned} \lambda_{\min} \left(\frac{1}{2}(\underline{B}^\iota(j) + \bar{B}^\iota(j)) \right) - \rho \left(\frac{1}{2}(\bar{B}^\iota(j) - \underline{B}^\iota(j)) \right) &\equiv c_B^\iota(j) \\ \text{for } j = 1, \dots, r^\iota, \quad \iota = 1, \dots, s, \end{aligned} \quad (3.5)$$

where λ_{\min} and ρ denotes the minimal eigenvalue and the spectral radius, respectively,

$$\begin{aligned} U_{ad}^{\kappa^\iota} &= \{2 \times 2 \text{ symmetric matrices } \kappa^\iota : \underline{\kappa}_{ik}^\iota(j) \leq \kappa_{ik}^\iota|_{\Omega_j^\iota} = \text{const.} \leq \bar{\kappa}_{ik}^\iota(j), \\ &j \leq r^\iota, \quad i, k \leq 2\} \end{aligned} \quad (3.6)$$

where $\underline{\kappa}^\iota(j)$ and $\bar{\kappa}^\iota(j)$ are given 2×2 symmetric matrices, $j = 1, \dots, r^\iota, \iota = 1, \dots, s$, and let there exist positive constants $c_B^\iota(j)$ such that

$$\lambda_{\min} \left(\frac{1}{2}(\underline{\kappa}^t(j) + \overline{\kappa}^t(j)) \right) - \rho \left(\frac{1}{2}(\overline{\kappa}^t(j) - \underline{\kappa}^t(j)) \right) \equiv c_{\kappa}^t(j) \text{ for } j \leq r^t, \ t \leq s, \quad (3.7)$$

where λ_{\min} and ρ denotes the minimal eigenvalue and the spectral radius, respectively. If (3.4) and (3.5) are satisfied, then the matrices $B^t(j) \equiv B_{|\Omega_j^t}^{B^t}$ are positive definite for any $B^t \in U_{ad}^{B^t}$, $t = 1, \dots, s$ and any $j \leq r^t$ (see [8]) and the matrices $\kappa^t(j) = \kappa_{|\Omega_j^t}^t$ are positive definite for any $\kappa^t \in U_{ad}^{\kappa^t}$, $t \leq s$, $j \leq r^t$.

Furthermore, we define

$$U_{ad}^{F_i^t} = \{f \in L^\infty(\Omega) : \underline{F}_i^t(j) \leq f_{|\Omega_j^t} = \text{const.} \leq \overline{F}_i^t(j), \ j \leq r^t\}, \quad (3.8)$$

for $i \leq 2$, $t \leq s$, where $\underline{F}_i^t(j)$ and $\overline{F}_i^t(j)$ are given constants;

$$U_{ad}^{W^t} = \{w \in L^\infty(\Omega) : \underline{W}^t(j) \leq w_{|\Omega_j^t} = \text{const.} \leq \overline{W}^t(j), \ j \leq r^t\}, \quad (3.9)$$

for $t \leq s$, where $\underline{W}^t(j)$ and $\overline{W}^t(j)$ are given constants;

$$U_{ad}^{T_1} = \{\mathcal{T} \in L^\infty(\Gamma_\tau) : \underline{T}_1(t) \leq \mathcal{T}_{|\Gamma_\tau^t} = \text{const.} \leq \overline{T}_1(t), \ t \leq s\}, \quad (3.10)$$

where $\underline{T}_1(t)$ and $\overline{T}_1(t)$ are given constants;

$$U_{ad}^{u_{0i}} = \{u \in L^\infty(\Gamma_u) : \underline{u}_{0i}(t) \leq u_{|\Gamma_u^t} = \text{const.} \leq \overline{u}_{0i}(t), \ t \leq s\}, \quad (3.11)$$

where $\underline{u}_{0i}(t)$ and $\overline{u}_{0i}(t)$, $i = 1, 2$, are given constants;

$$U_{ad}^{P_i} = \{P \in L^\infty(\Gamma_\tau) : \underline{P}_i(t) \leq p_{|\Gamma_\tau^t} = \text{const.} \leq \overline{P}_i(t), \ t \leq s\}, \quad (3.12)$$

where $\underline{P}_i(t)$ and $\overline{P}_i(t)$, $i = 1, 2$, are given constants;

$$U_{ad}^{\beta_i^t} = \{b \in L^\infty(\Omega) : \underline{\beta}_i^t(j) \leq b_{|\Omega_j^t} = \text{const.} \leq \overline{\beta}_i^t(j), \ j \leq r^t\}, \quad (3.13)$$

for $i \leq 3$, $t \leq s$, where $\underline{\beta}_i^t(j)$ and $\overline{\beta}_i^t(j)$ are given constants.

Finally, we define the set of admissible data by

$$\begin{aligned} U_{ad} = & \quad \prod_{t \leq s} U_{ad}^{B^t} \times \prod_{t \leq s} U_{ad}^{\kappa^t} \times \prod_{t \leq s, i \leq 2} U_{ad}^{F_i^t} \times \prod_{t \leq s} U_{ad}^{W^t} \times \\ & \quad \times \prod_{t \leq s, i \leq 2} U_{ad}^{\beta_i^t} \times \prod_{i \leq 2} U_{ad}^{P_i} \times \prod_{i \leq 2} U_{ad}^{u_{0i}} \times \prod_{t \leq s} U_{ad}^{T_1}. \end{aligned} \quad (3.14)$$

Further, instead of $b(T, z)$, $a(\mathbf{u}, \mathbf{v})$, $s(z)$, $S(\mathbf{v}, T)$ we will write $b(A; T, z)$, $a(A; \mathbf{u}, \mathbf{v})$, $s(A; z)$, $S(A; \mathbf{v}, T)$ for any $A \in U_{ad}$.

The next results are parallel to those of [3] for the general case with friction.

Lemma 1. There exist positive constants c_i , $i = 0, 1, \dots, 5$ independent of $A \in U_{ad}$, such that

$$b(A; z, z) \geq c_0 \|z\|_{W_1}^2 \quad \forall z \in V_1, \quad (3.15)$$

$$|b(A; z, y)| \leq c_1 \|z\|_{W_1} \|y\|_{W_1} \quad \forall z, y \in W_1, \quad (3.16)$$

$$a(A; \mathbf{v}, \mathbf{v}) \geq c_2 \|\mathbf{v}\|_W^2 \quad \forall \mathbf{v} \in V, \quad (3.17)$$

$$|a(A; \mathbf{v}, \mathbf{w})| \leq c_3 \|\mathbf{v}\|_W \|\mathbf{w}\|_W \quad \forall \mathbf{v}, \mathbf{w} \in W, \quad (3.18)$$

$$|s(A; z)| \leq c_4 \|z\|_{0, \Omega} \quad \forall z \in V_1, \quad (3.19)$$

$$|S(A; \mathbf{v}, T)| \leq c_5 (\|\mathbf{v}\|_{0, \Omega} + \|\mathbf{v}\|_{0, \Gamma_\tau} + \|T - T_0\|_{0, \Omega} \|\mathbf{v}\|_W) \quad \forall \mathbf{v}, \mathbf{w} \in W. \quad (3.20)$$

Proposition 1. There exists a unique weak solution $(T(A), \mathbf{u}(A))$ of the problem (\mathcal{P}) for any $A \in U_{ad}$. Moreover, $\|T(A)\|_{W_1} \leq c$, where c is independent of A .

To find the most “dangerous” input data A in the set U_{ad} , we will introduce a criterion, i.e. defined a functional, which depends on the solution $(T(A), \mathbf{u}(A))$ of the problem (\mathcal{P}) . Such criteria can be as follows:

Let $G_r \subset \bigcup_{t \leq s} \Omega^t$, $r = 1, \dots, \bar{r}$, be subdomains adjacent to the boundaries $\partial\Omega^t$. Then we define

$$\Phi_1(T) = \max_{r \leq \bar{r}} \varphi_r(T) = \max_{r \leq \bar{r}} \left[(\text{meas}_2 G_r)^{-1} \int_{G_r} T d\mathbf{x} \right]; \quad (3.21)$$

let $G'_r \subset \Gamma_u$, $r \leq \bar{r}$ and

$$\Phi_2(T) = \max_{r \leq \bar{r}} \psi_r(T) = \max_{r \leq \bar{r}} \left[(\text{meas}_1 G'_r)^{-1} \int_{G'_r} T ds \right]; \quad (3.22)$$

and

$$\Phi_3(\mathbf{u}) = \max_{r \leq \bar{r}} \chi_r(\mathbf{u}) = \max_{r \leq \bar{r}} \left[(\text{meas}_2 G_r)^{-1} \int_{G_r} u_i n_i(X_r) d\mathbf{x} \right]; \quad (3.23)$$

where $\mathbf{n}(X_r)$ is the unit outward normal at a fixed point $X_r \in \partial\Omega^t \cap \partial G_r$ (if $G_r \subset \Omega^t$) to the boundary $\partial\Omega^t$;

$$\Phi_4(\mathbf{u}) = \max_{r \leq \bar{r}} \chi'_r(\mathbf{u}) = \max_{r \leq \bar{r}} \left[(\text{meas}_1 G'_r)^{-1} \int_{G'_r} u_i n_i(X_r) ds \right]; \quad (3.24)$$

where $G'_r = \bigcup_{t \leq s} \partial\Omega^t \setminus \Gamma_u$. Since the weak solution $\mathbf{u}(A)$ of our problem (2.10) depends on $T(A)$, then $\mathbf{u}(A) = \mathbf{u}(A; T(A))$ and instead of $\Phi_i(\mathbf{u})$ we write $\Phi_i(A; \mathbf{u}, T)$. Thus we may define

$$\Phi_5(A; u, T) = \max_{r \leq \bar{r}} \omega_r(A; u, T) = \max_{r \leq \bar{r}} \left[(\text{meas}_2 G_r)^{-1} \int_{G_r} I_2^2(\tau(A; \mathbf{u}, T)) d\mathbf{x} \right]; \quad (3.25)$$

here $I_2(\tau) = \left(\sum_{i,j=1}^3 \tau_{ij}^D \tau_{ij}^D \right)^{\frac{1}{2}}$ is the intensity of shear stress, where $\tau_{ij}^D = \tau_{ij} - \frac{1}{3} \tau_{kk} \delta_{ij}$ and $\tau(A; \mathbf{u}, T)$ is defined by (2.2). Finally, we may choose

$$\Phi_6(A; u, T) = \max_{r \leq \bar{r}} \mu_r(A; u, T) = \max_{r \leq \bar{r}} \left[(\text{meas}_2 G_r)^{-1} \int_{G_r} (-\tau_n(A; \mathbf{u}, T)) ds \right]; \quad (3.26)$$

where G_r is a small subdomain adjacent to Γ_c .

Now we formulate the worst scenario problems as follows:
find

$$A^{0i} = \arg \max_{A \in U_{ad}} \Phi_i(T(A)), \quad i = 1, 2 \quad (3.27)$$

and

$$A^{0i} = \arg \max_{A \in U_{ad}} \Phi_i(\mathbf{u}(A), T(A)), \quad i = 3, 4, 5, 6, \quad (3.28)$$

where $(T(A), \mathbf{u}(A))$ is weak solution of the problem (\mathcal{P}) .

4 Stability of Weak Solutions

To prove the solvability of worst scenario problems (3.27), (3.28), we have to study the mapping $A \mapsto T(A)$, $A \mapsto \mathbf{u}(A, T(A))$. We introduce the decomposition of $A \in U_{ad}$ as $A = \{A', A''\}$, where

$$A' = \{\Pi_{l \leq s} \Pi_{j \leq r^l} \kappa^l(j), \Pi_{l \leq s} \Pi_{j \leq r^l} W^l(j), \Pi_{l \leq s} T_1^l\}, \quad A' \in R^{p_1}, \quad p_1 = 4 \sum_{l \leq s} r^l + s,$$

and

$$\begin{aligned} A'' &= \{\Pi_{l \leq s} \Pi_{j \leq r^l} B^l(j), \Pi_{l \leq s} \Pi_{j \leq r^l} \mathbf{F}^l(j), \Pi_{l \leq s} \mathbf{P}^l, \Pi_{l \leq s} \mathbf{u}_0^l, \Pi_{l \leq s} \Pi_{j \leq r^l} \beta^l(j)\}, \\ A'' &\in R^{p_2}, \quad p_2 = \left(\sum_{l \leq s} r^l \right) [9 + 2(1 + 2s)]. \end{aligned}$$

We are going to show the continuity of the mappings $A' \mapsto T(A')$, $A \mapsto \mathbf{u}(A, T(A'))$ for $A' \in U'_{ad} = \Pi_{l \leq s} U_{ad}^{\kappa^l} \times \Pi_{l \leq s} U_{ad}^{W^l} \times U_{ad}^{T_1^l}$ and $A'' \in U''_{ad} = \Pi_{l \leq s} U_{ad}^{B^l} \times \Pi_{l \leq s, i \leq 2} U_{ad}^{F_i^l} \times \Pi_{l \leq s, i \leq 2} U_{ad}^{\beta_i^l} \times \Pi_{l \leq 2} U_{ad}^{P_i} \times \Pi_{i \leq 2} U_{ad}^{u_{0i}}$, respectively. Since the problem discussed is quasi-coupled, we will prove the following theorems and lemma:

Theorem 1. Let $A' \in U'_{ad}$, $A'_n \rightarrow A'$ in R^{p_1} as $n \rightarrow \infty$. Then

$$T(A'_n) \rightarrow T(A) \quad \text{in } W_1.$$

Sketch of the proof: Since

$$b(A; z, z) \geq \left(\min_{l \leq s, j \leq r^l} c_\kappa^l(j) \right) \sum_{l \leq s} \int_{\Omega^l} |\text{grad } z^l|^2 d\mathbf{x}, \quad (4.1)$$

for $T_n := T(A'_n)$ we obtain $\|T_n\|_{W_1} \leq c$ for all n . Then a $T \in W_1$ and a subsequence $\{T_m\} \subset \{T_n\}$ exist such that

$$T_m \rightharpoonup T \quad \text{weakly in } W_1. \quad (4.2)$$

By definition

$$b(A'_m; T_m, z) = s(A'_m; z) \quad \forall z \in V_1, \quad \forall m. \quad (4.3)$$

Since

$$\begin{aligned} |b(A'_m; T_m, z) - b(A'; T, z)| &\rightarrow 0, \quad \text{as } m \rightarrow \infty, \\ |s(A'_m; z) - b(A'; z)| &\rightarrow 0, \quad \text{as } m \rightarrow \infty, \end{aligned}$$

we prove that

$$b(A'_m; T_m, z) \rightarrow b(A'; T, z) \quad \text{as } m \rightarrow \infty, \quad (4.4)$$

$$s(A'_m; z) \rightarrow s(A'; z) \quad \text{as } m \rightarrow \infty. \quad (4.5)$$

Then we pass to the limit with $m \rightarrow \infty$ in (4.3). Using (4.4), (4.5) we prove that $T = T(A')$ is a weak solution of thermal part of the problem. Since it is unique, the whole sequence $\{T_n\}$ tends $T(A')$ weakly in W_1 . \square

Remark 2. It can be proved that $T_m \rightarrow T$ converges also strongly in W_1 .

Lemma 2. If $A_n'' \in U_{ad}$, $A_n'' \rightarrow A''$ in R^{p_2} , and $\mathbf{u}_n \rightarrow u$ weakly in W , then

$$a(A_n''; \mathbf{u}_n, \mathbf{v}) \rightarrow a(A''; \mathbf{u}, \mathbf{v}) \quad \forall \mathbf{v} \in W, \quad (4.6)$$

$$S(A_n''; \mathbf{u}_n, T) \rightarrow S(A''; \mathbf{u}, T) \quad \forall T \in W_1. \quad (4.7)$$

Sketch of the proof: The proof follows from the fact that

$$\begin{aligned} |a(A_n''; \mathbf{u}_n, \mathbf{v}) - a(A''; \mathbf{u}, \mathbf{v})| &\rightarrow 0 \quad \text{for } n \rightarrow \infty, \\ |S(A_n''; \mathbf{u}_n, T) - S(A''; \mathbf{u}, T)| &\rightarrow 0 \quad \text{for } n \rightarrow \infty. \end{aligned}$$

□

Theorem 2. Let $A_n \in U_{ad}$, $A_n \rightarrow A$ in $U \equiv R^{p_2}$. Then

$$\mathbf{u}(A_n) \rightarrow \mathbf{u}(A) \quad \text{in } W. \quad (4.8)$$

Sketch of the proof: Let us denote $\mathbf{u}_n := \mathbf{u}(A_n)$, $\mathbf{u} := \mathbf{u}(A)$, $\mathbf{u}_{0n} := \mathbf{u}_0(A_n)$, $\mathbf{u}_0 := \mathbf{u}_0(A)$, $T_n := T(A_n)$, $T := T(A)$. Inserting $\mathbf{u} := \mathbf{u}_0 + \mathbf{w}(A)$, $\mathbf{w}(A) \in K$, $\mathbf{u}_n := \mathbf{u}_{0n} + \mathbf{w}_n(A)$, $\mathbf{w}_n(A) \in K$, $\mathbf{v} := \mathbf{u}_0 + \mathbf{w}$ or $\mathbf{v} := \mathbf{u}_{0n} + \mathbf{w}$, $\mathbf{w} \in K$ into the variational inequality (2.10), we obtain

$$a(A_n; \mathbf{w}_n, \mathbf{w} - \mathbf{w}_n) \geq S(A_n; \mathbf{w} - \mathbf{w}_n, T_n) - a(A_n; \mathbf{u}_{0n}, \mathbf{w} - \mathbf{w}_n). \quad (4.9)$$

Hence, putting $\mathbf{w} = 0$, using Lemma 1, Theorem 1, definition of $U_{ad}^{u_0}$, after some modifications we find that

$$c_0 \|w_n\|_W^2 \leq c_7 \|w_n\|_W + c_8.$$

As a consequence, \mathbf{w}_n are bounded in W and there exists a subsequence $\{\mathbf{w}_k\}$ and a function $\boldsymbol{\omega} \in W$ such that

$$\mathbf{w}_k \rightarrow \boldsymbol{\omega} \quad \text{weakly in } W, \text{ as } k \rightarrow \infty. \quad (4.10)$$

It can be shown that $\boldsymbol{\omega} = \mathbf{w}(A)$. Thus, since $\boldsymbol{\omega} \in K$ and since $a(A_k; \mathbf{w}_k - \boldsymbol{\omega}, w_k - \boldsymbol{\omega}) \geq 0$, after some modification and using Lemma 2, we obtain $\liminf a(A_k; \mathbf{w}_k, \mathbf{w}_k - \boldsymbol{\omega}) \geq \lim a(A_k; \boldsymbol{\omega}, w_k - \boldsymbol{\omega}) = 0$. Inserting $\mathbf{w} := \boldsymbol{\omega}$ into (4.9) we arrive at

$$a(A_k; w_k, \boldsymbol{\omega} - w_k) \geq S(A_k; \boldsymbol{\omega} - w_k, T_k) - a(A_k; \mathbf{u}_{0k}, \boldsymbol{\omega} - w_k)$$

and

$$\limsup a(A_k; \mathbf{w}_k, \mathbf{w}_k - \boldsymbol{\omega}) \leq \limsup S(A_k; \mathbf{w}_k - \boldsymbol{\omega}, T_k) + \limsup a(A_k; \mathbf{u}_{0k}, \boldsymbol{\omega} - \mathbf{w}_k).$$

For any $A \in U_{ad}$, $T \in W_1$ we can show that $\lim S(A_k; \mathbf{w}_k - \boldsymbol{\omega}, T_k) = 0$ and $\lim a(A_k; \mathbf{w}_k, \mathbf{w}_k - \boldsymbol{\omega}) = 0$ as $\limsup a(A_k; \mathbf{w}_k, \mathbf{w}_k - \boldsymbol{\omega}) \leq 0$, from which it follows that $\lim a(A_k; \mathbf{w}_k, \mathbf{w}_k - \boldsymbol{\omega}) = 0$. It can be shown that $|a(A_k; \mathbf{w}_k, \mathbf{w} - \mathbf{w}_k) - a(A; \boldsymbol{\omega}, \mathbf{w} - \boldsymbol{\omega})| \rightarrow 0$; then

$$\lim a(A_k; w_k, w - w_k) = a(A; \boldsymbol{\omega}, w - \boldsymbol{\omega})$$

and since $|S(A_k; \mathbf{w} - \mathbf{w}_k, T_k) - S(A; \mathbf{w} - \boldsymbol{\omega}, T)| \rightarrow 0$, then

$$\lim S(A_k; w - w_k, T_k) = S(A; w - \boldsymbol{\omega}, T).$$

Moreover, we have $|a(A_k; \mathbf{w} - \mathbf{w}_k, \mathbf{u}_{0k}) - a(A; \mathbf{w} - \boldsymbol{\omega}, \mathbf{u}_{0k})| \rightarrow 0$, where Lemma 1, Lemma 2 and the convergence $\mathbf{u}_{0k} \rightarrow \mathbf{u}_0$ in W were used. Thus

$$\lim a(A_k; \mathbf{w} - \mathbf{w}_k, \mathbf{u}_{0k}) = a(A; \mathbf{w} - \boldsymbol{\omega}, \mathbf{u}_0).$$

Passing to the limit with $k \rightarrow \infty$, we obtain

$$a(A; \boldsymbol{\omega}, \mathbf{w} - \boldsymbol{\omega}) \geq S(A; \mathbf{w} - \boldsymbol{\omega}, T) - a(A; \mathbf{w} - \boldsymbol{\omega}, u_0). \quad (4.11)$$

Since the variational inequality (2.10) has a unique solution, $\boldsymbol{\omega} = \mathbf{w}(A)$ follows from (4.11) and moreover, whole sequence $\{\mathbf{w}(A_n)\}$ tends to $\mathbf{w}(A)$ weakly in W .

Furthermore, the strong convergence can also be proved.

5 Existence of a Solution of the Worst Scenario Problem

To prove the existence of a solution of the worst scenario problem, we will use the following lemma.

Lemma 3.

(i) Let $\Phi_i(T)$, $i = 1, 2$, be defined by (3.21), (3.22) and let $T_n \rightarrow T$ in W_1 , as $n \rightarrow \infty$. Then

$$\lim_{n \rightarrow \infty} \Phi_i(T_n) = \Phi_i(T), \quad i = 1, 2. \quad (5.1)$$

(ii) Let $\Phi_i(\mathbf{u})$, $i = 3, 4$, be defined by (3.23), (3.24) and let $\mathbf{u}_n \rightarrow \mathbf{u}$ in W , as $n \rightarrow \infty$. Then

$$\lim_{n \rightarrow \infty} \Phi_i(\mathbf{u}_n) = \Phi_i(\mathbf{u}), \quad i = 3, 4. \quad (5.2)$$

(iii) Let $\Phi_i(A; \mathbf{u}, \mathbf{T})$, $i = 5, 6$, be defined by (3.25), (3.26) and let $A_n \rightarrow A$ in U , $A_n \in U_{ad}$, $\mathbf{u}_n \rightarrow \mathbf{u}$ in W and $T_n \rightarrow T$ in $L^2(\Omega)$, as $n \rightarrow \infty$. Then

$$\lim_{n \rightarrow \infty} \Phi_i(A_n, \mathbf{u}_n, T_n) = \Phi_i(A, \mathbf{u}, T), \quad i = 5, 6 \quad (5.3)$$

The proof is a modification of that of [3].

As the main result of the paper we present the following theorem:

Theorem 3. There exists at least one solution of the worst scenario problems (3.27), (3.28), $i = 1, \dots, 6$. The proof is a modification of that of [3].

6 Conclusion

Mathematical models connected with the safety of construction and of operation of the radioactive waste repositories involve input data (thermal conductivity and elastic coefficients, body and surface forces, thermal sources, coefficients of thermal expansion, boundary values, coefficient of friction on contact boundaries, etc.) which cannot be determined uniquely, but only in some intervals, given by the accuracy of measurements and the approximate solutions of identification problems. The notation "reliable solution" denotes the worst case among a set of possible solutions where the degree of badness is measured by a criterion functional. For the safety of the radioactive waste repositories we seek the maximal value of this functional, which depends on the solution of the mathematical model. Then for the computations of such problems (some mean values of temperatures, displacements, intensity of shear stresses, principal stresses, stress tensor components, normal and tangential components of the displacement or stress vector on the contact boundaries, etc.) we have to formulate a corresponding maximization (worst scenario) problem. Then methods and algorithms known from "optimal design" can be used.

To construct a model of structures under the influence of critical conditions the influence of global tectonics onto a local area, where the critical structure is built as well as the influence of the resulting local geomechanical processes on a critical structure must be taken into account ([6]). Problems of this kind with uncertain input data are problems with high level radioactive waste repositories. In the

case of the high level radioactive waste repositories the effects of geodynamical processes in the sense of plate tectonics must be taken into consideration, namely in regions near tectonic areas (e.g. the Japan island arc, the Central and South Europe, etc), but also in the platform regions (as in Sweden, Canada, etc.). Another example is represented by modelling an interaction between a tunnel wall and a rock massif in the radioactive waste repository tunnels or by modelling of a tunnel crossing by an active deep fault(s), respectively.

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